## On the Connection between Generalized Hypergeometric Functions and Dilogarithms

## M.A. Sanchis-Lozano

Departamento de Física Teórica and Instituto de Física Corpuscular ( IFIC ) Centro Mixto Universidad de Valencia-CSIC

Dr. Moliner 50, 46100 Burjassot, Valencia (Spain) February 7, 2008

#### Abstract

Several integrals involving powers and ordinary hypergeometric functions are rederived by means of a generalized hypergeometric function of two variables (Appell's function) recovering some well-known expressions as particular cases. Simple connections between dilogarithms and a kind of Appell's function are shown. A relationship is generalized to polylogarithms.

FTUV: 95/49 IFIC: 95/51 Hypergeometric functions play an important role in mathematical physics since they are related to a wide class of special functions appearing in a large variety of fields. In particular, it is well-known a long time ago that integrals emerging from loop calculations in Feynman diagrams can be written in terms of hypergeometric functions [1]. More recently, generalized hypergeometric functions of one or several variables have been used in the evaluation of scalar one-loop Feynman integrals [2] or multiloop ones [3].

In this work we firstly rederive <sup>1</sup> an integral expression involving two ordinary Gauss' functions yielding a generalized hypergeometric function of two variables (Appell's function). Several formulae appearing in standard tables (e.g. Gradshteyn and Ryzhik [5]) of utility for the evaluation of Feynman loop integrals are obtained as particular cases. Moreover, we have shown a simple relationship between a kind of Appell's function and dilogarithms [6], contributing to enlarge the knowledge on the connection between them. In the appendices at the end of the paper we present a brief survey on the generalized Gauss' functions establishing the notation employed and revising some of their properties needed in this work.

#### expression 1

$$\int_{0}^{1} du \ u^{\gamma-1} (1-u)^{\rho-1} {}_{2}F_{1}[\sigma, \eta; \gamma; zu] {}_{2}F_{1}[\alpha, \beta; \rho; k(1-u)] =$$

$$= \frac{\Gamma(\gamma)\Gamma(\rho)}{\Gamma(\gamma+\rho)} F_{3}[\alpha, \sigma, \beta, \eta; \gamma+\rho; k, z]$$
(1)

provided that  $Re(\gamma) > 0$ ,  $Re(\rho) > 0$ ,  $|arg(1-k)| < \pi$ ,  $|arg(1-z)| < \pi$ .

*Proof.* We will first show that Eq. (1) holds in the domain of convergence of the series. Expanding one of the two  ${}_{2}F_{1}$  functions as a power series leads to:

$$\sum_{n=0}^{\infty} \frac{(\alpha)_n(\beta)_n}{(\rho)_n} \frac{k^n}{n!} \int_0^1 du \ u^{\gamma-1} (1-u)^{\rho+n-1} \ _2F_1[\sigma, \eta; \gamma; zu] \quad ; \ |k| < 1, \ |z| < 1$$

where we have interchanged the order of summation and integration on account of the dominated convergence theorem of Lebesgue, provided that  $Re(\gamma) > 0$ ,  $Re(\rho) > 0$ . Now, performing the integration over u one gets from (A.3):

$$\Gamma(\gamma) \sum_{n=0}^{\infty} \frac{(\alpha)_n(\beta)_n}{(\rho)_n} \frac{k^n}{n!} \frac{\Gamma(\rho+n)}{\Gamma(\gamma+\rho+n)} {}_{3}F_{2}[\gamma, \sigma, \eta; \gamma+\rho+n, \gamma; z] =$$

$$= \Gamma(\gamma) \sum_{n=0}^{\infty} \frac{(\alpha)_n(\beta)_n}{(\rho)_n} \frac{k^n}{n!} \frac{\Gamma(\rho+n)}{\Gamma(\gamma+\rho+n)} {}_{2}F_{1}[\sigma, \eta; \gamma+\rho+n; z]$$

where a cancellation between two parameters in the  $_3F_2$  function occurred.

<sup>&</sup>lt;sup>1</sup>An exhaustive set of integrals involving generalized Gauss functions containing ours as a particular case can be found in [4].

Finally, using that:  $\Gamma(\rho+n) = \Gamma(\rho)(\rho)_n$ ,  $\Gamma(\gamma+\rho+n) = \Gamma(\gamma+\rho)(\gamma+\rho)_n$ , one arrives at

$$\frac{\Gamma(\gamma)\Gamma(\rho)}{\Gamma(\gamma+\rho)} \sum_{n=0}^{\infty} \frac{(\alpha)_n(\beta)_n}{(\gamma+\rho)_n} \frac{k^n}{n!} {}_{2}F_{1}[\sigma,\eta;\gamma+\rho+n;z]$$

which leads to Eq. (1) at once in virtue of (B.2). Moreover, the integral (1) furnishes a single-valued function of two variables beyond the domain of convergence of the series by imposing the cuts:  $|arg(1-k)| < \pi$ ,  $|arg(1-z)| < \pi$ .

#### expression 1.1

Setting  $\rho = \delta - \gamma$ ,  $\alpha = \delta - \sigma$  and  $\beta = \delta - \eta$  in Eq. (1) the formula 7.512.7 of reference [5] is recovered: <sup>2</sup>

$$\int_0^1 du \ u^{\gamma-1} (1-u)^{\delta-\gamma-1} \ _2F_1[\sigma, \eta; \gamma; zu] \ _2F_1[\delta - \sigma, \delta - \eta; \delta - \gamma; k(1-u)] =$$

$$= \frac{\Gamma(\gamma)\Gamma(\delta - \gamma)}{\Gamma(\delta)} (1-k)^{\sigma+\eta-\delta} \ _2F_1[\sigma, \eta; \delta; k + z - kz] \tag{2}$$

provided that  $Re(\delta) > Re(\gamma) > 0$ ,  $|arg(1-k)| < \pi$ ,  $|arg(1-z)| < \pi$ .

This can be directly obtained from Eq. (1) taking further into account the property (B.6) which here implies:

$$F_3[\delta - \sigma, \sigma, \delta - \eta, \eta; \delta; k, z] = (1 - k)^{\sigma + \eta - \delta} {}_{2}F_1[\sigma, \eta; \delta; k + z - kz]$$

#### expression 1.2

With the aid of (A.4) the left hand side of Eq. (1) can be written as:

$$\int_0^1 du \ u^{\gamma-1} (1-u)^{\rho-1} (1-k(1-u))^{-\alpha} \ _2F_1[\sigma,\eta;\gamma;zu] \ _2F_1\left[\alpha,\rho-\beta;\rho;\frac{k(1-u)}{(k-1-ku)}\right]$$

Now, let us assume that z and k are related through k = z/(z-1). Then

$$(1-z)^{\alpha} \int_{0}^{1} du \, u^{\gamma-1} (1-u)^{\rho-1} (1-zu)^{-\alpha} \, {}_{2}F_{1}[\sigma, \eta; \gamma; zu] \, {}_{2}F_{1}[\alpha, \rho-\beta; \rho; \frac{z(1-u)}{1-zu}] =$$

$$= \frac{\Gamma(\gamma)\Gamma(\rho)}{\Gamma(\gamma+\rho)} \, F_{3}[\alpha, \sigma, \beta, \eta; \gamma+\rho; z/(z-1), z]$$

Next, let us suppose further that  $\beta = \gamma + \rho - \eta$ . Then taking into account consecutively the properties (B.5) and (B.4) the right hand side of the last expression becomes:

$$\frac{\Gamma(\gamma)\Gamma(\rho)}{\Gamma(\gamma+\rho)} (1-z)^{\alpha} F_1[\eta;\sigma;\alpha;\gamma+\rho;z,z] = \frac{\Gamma(\gamma)\Gamma(\rho)}{\Gamma(\gamma+\rho)} (1-z)^{\alpha} {}_2F_1[\sigma+\alpha,\eta;\gamma+\rho;z]$$

<sup>&</sup>lt;sup>2</sup>Except the exponent of (1-k) which in our notation would read:  $2\sigma - \delta$ . Clearly this is an error since the result should be invariant under the interchange of  $\sigma$  and  $\eta$ , as the l.h.s. certainly is. The original source [7] is equally wrong.

Hence one recovers the formula 7.512.8 of ref. [5]:

$$\int_{0}^{1} du \ u^{\gamma-1} (1-u)^{\rho-1} (1-zu)^{-\alpha} \ _{2}F_{1}[\sigma, \eta; \gamma; zu] \ _{2}F_{1}\left[\alpha, \eta - \gamma; \rho; \frac{z(1-u)}{(1-zu)}\right] =$$

$$= \frac{\Gamma(\gamma)\Gamma(\rho)}{\Gamma(\gamma+\rho)} \ _{2}F_{1}[\sigma + \alpha, \eta; \gamma + \rho; z]$$
(3)

provided that  $Re(\gamma) > 0$ ,  $Re(\rho) > 0$ ,  $|arg(1-z)| < \pi$ .

Let us now go back again to Eq. (1) and consider  $\eta = \gamma$  as a new special case. Then two parameters of a hypergeometric function in the integrand cancel, i.e.  ${}_2F_1[\sigma, \gamma; \gamma; zu] = {}_1F_0[\sigma; zu] = (1 - zu)^{-\sigma}$ , yielding:

#### expression 2

$$\int_0^1 du \ u^{\gamma-1} (1-u)^{\rho-1} (1-zu)^{-\sigma} \ _2F_1[\alpha,\beta;\gamma;ku] =$$

$$= \frac{\Gamma(\gamma)\Gamma(\rho)}{\Gamma(\gamma+\rho)} (1-z)^{-\sigma} F_3[\alpha,\sigma,\beta,\rho;\gamma+\rho;k,z/(z-1)] \tag{4}$$

provided that  $Re(\gamma) > 0$ ,  $Re(\rho) > 0$ ,  $|arg(1-k)| < \pi$ ,  $|arg(1-z)| < \pi$ .

*Proof.* It follows inmediately as a particular case of the expression **1** by means of the change of the integration variable:  $u\rightarrow 1-u$  and interchanging the  $\gamma$  and  $\rho$  parameters.

An alternative (direct) proof is achieved with the aid of the integral representation of the  $F_3$  Appell's function. Starting from (B.3) and making the consecutive changes of the integration variables:  $v \rightarrow 1-v$  and  $u \rightarrow uv$  it follows that

$$\frac{\Gamma(\beta)\Gamma(\rho)\Gamma(\gamma-\beta)}{\Gamma(\gamma+\rho)} F_3[\alpha,\sigma,\beta,\rho;\gamma+\rho;k,z/(z-1)] = (1-z)^{\sigma} \int_0^1 \int_0^1 dv \ du \ v^{\gamma-1} u^{\beta-1} (1-v)^{\rho-1} (1-u)^{\gamma-\beta-1} (1-zv)^{-\sigma} (1-kvu)^{-\alpha}$$

Hence the expression **2** is immediately obtained by expressing  ${}_{2}F_{1}[\alpha, \beta; \gamma; kv]$  in its Euler's integral representation (A.2).

#### expression 2.1

Setting k = 1 in expression 2 reproduces the result 7.512.9 of ref. [5]:

$$\int_0^1 du \ u^{\gamma-1} (1-u)^{\rho-1} (1-zu)^{-\sigma} \ _2F_1[\alpha,\beta;\gamma;u] =$$

$$= \frac{\Gamma(\gamma)\Gamma(\rho)\Gamma(\gamma+\rho-\alpha-\beta)}{\Gamma(\gamma+\rho-\alpha)\Gamma(\gamma+\rho-\beta)} (1-z)^{-\sigma} \ _3F_2[\rho,\sigma,\gamma+\rho-\alpha-\beta;\gamma+\rho-\alpha,\gamma+\rho-\beta;z/(z-1)]$$
provided additionally that:  $Re(\gamma+\rho-\alpha-\beta) > 0$ .

This can be easily shown by rewriting the power expansion of  $F_3$  following (B.2) in terms of  ${}_2F_1[\alpha,\beta;\gamma+\rho+n;1]$  supposed the convergence of the series, and using the Gauss' summation relation:

$$_{2}F_{1}[a,b;c;1] = \frac{\Gamma(c)\Gamma(c-a-b)}{\Gamma(c-a)\Gamma(c-b)}$$
;  $Re(c-a-b) > 0$ 

#### Particular values of the parameters of $F_3$

Let us now take the particular values of the parameters:  $\alpha = \beta = \sigma = \eta = \gamma = 1$  and  $\rho = 2$  in the  $F_3[\alpha, \sigma, \beta, \eta; \gamma + \rho; x, y]$  Appell's function. From our expression 1, it is easy to see that making y = 1 ones gets

$$x F_3[1,1,1,1;3;x,y=1] = 2x {}_{3}F_2[1,1,1;2,2;x] = 2 Li_2(x)$$
 (5)

where Eq. (A.8) has been taken into account. In fact, one can get the same result by expanding  $F_3[1,1,1,1;3;x,1]$  in terms of  ${}_2F_1[1,1;3+m;1]$  and using the Gauss' summation relation. (See also appendix B for a generalization to polylogarithms.) The restriction |x| < 1 can be dropped on account of analytic continuation, extending the domain of analyticity over the complex x-plane cut from 1 to  $\infty$  along the real axis. It is obvious from symmetry, that an equivalent expression for y must be satisfied. In fact, Eq. (5) can be viewed as a particular case of a more general relationship between this Appell's series and dilogarithms:

#### expression 3

$$\frac{1}{2}xy \ F_3[1,1,1,1;3;x,y] = Li_2(x) + Li_2(y) - Li_2(x+y-xy)$$

$$|arg(1-x)| < \pi, \ |arg(1-y)| < \pi.$$
(6)

*Proof.* This formula can be again derived from expression  $\mathbf{1}$  by calculating directly the integral in terms of dilogarithms. Instead, we will prove it by differentiating both sides with respect to x and y consecutively. Expanding the Appell's function as a double series, the result of differentiating the l.h.s. reads:

$$\frac{1}{2} F_3[1, 2, 2, 1; 3; x, y] \tag{7}$$

Now, invoking the property (B.6):

$$F_3[\alpha, \gamma - \alpha, \beta, \gamma - \beta; \gamma; x, y] = (1 - y)^{\alpha + \beta - \gamma} {}_2F_1[\alpha, \beta; \gamma; x + y - xy]$$

which is valid in a suitable small open polydisc centered at the origin, we conclude that (7) can be rewritten as

$$\frac{1}{2} {}_{2}F_{1}[1,2;3;x+y-xy]$$

Next, differentiating twice the r.h.s of Eq. (6) one gets:

$$-\frac{1}{x+y-xy}\left[1+\frac{1}{x+y-xy}\ln\left(1-(x+y-xy)\right)\right] = \frac{1}{2} {}_{2}F_{1}[1,2;3;x+y-xy]$$

the last step coming from (A.5). Then both sides in Eq. (6) would differ in f(x) + g(y):

$$\frac{1}{2}xy F_3[1,1,1,1;3;x,y] = Li_2(x) + Li_2(y) - Li_2(x+y-xy) + f(x) + g(y)$$

where f(x) and g(y) are functions to be determined by taking particular values of the variables. Setting x = 0 and y = 0 it is easy to see that  $f(x) = g(y) \equiv 0$ .

Now, by analytic continuation we dispense with the restriction on the small polydisc, extending its validity to a suitable domain of  $C^2$ : in order to get a single-valued function, with a well-defined branch for each dilogarithm in Eq. (6), we assume further that  $|arg(1-x)| < \pi$ ,  $|arg(1-y)| < \pi$ .

#### expression 3.1

$$x^{2} F_{3}[1, 1, 1, 1; 3; x, -x] = Li_{2}(x^{2})$$
 (8)

for x real.

*Proof.* It follows directly from Eq. (6) using the relation  $Li_2(x) + Li_2(-x) = \frac{1}{2}Li_2(x^2)$ .

#### expression 3.2

$$x^{2} F_{3}[1,1,1,1;3;x,x] = 4 Li_{2}\left(\frac{1}{2-x}\right) + 2 \ln^{2}(2-x) - \frac{\pi^{2}}{3}$$
 (9)

for x real and less than unity.

*Proof.* It follows directly from Eq. (6) using the relation:  $Li_2(2x - x^2) = 2Li_2(x) - 2Li_2(1/(2-x)) + \pi^2/6 - \ln^2(2-x)$  [6].

#### expression 3.3

$$\lim_{y\to 0} \frac{xy}{2} F_3[1,1,1,1;3;x,y] = y \left[ 1 + \frac{1-x}{x} \ln(1-x) \right]$$
 (10)

*Proof.* It follows directly from Eq. (6) using the relation:  ${}_2F_1[1,1;3;x] = (1-x)^{-1} {}_2F_1[1,2;3;x/(x-1)]$  and (A.5). If besides  $x \rightarrow 0$ , the limit xy/2 is quickly recovered.

The set of expressions 3 provide new connections (not shown in literature to our knowledge) between dilogarithms and a certain  $F_3$  Appell's function.

<sup>&</sup>lt;sup>3</sup>Observe that then each function of one complex variable obtained from (6) by fixing the other variable is analytic in the corresponding subset of  $C^2$ . Then the function of two variables  $\frac{1}{2}xyF_3$  is analytic according to the theorem of Hartogs-Osgood

## Appendices

## $\mathbf{A}$

Generalized Gauss' Functions of one variable

Hypergeometric functions can be introduced at first as series within a certain domain of convergence [8] [9] [10]. We write, using the abbreviate notation:

$$_{p}F_{q}[\{a\}_{p};\{b\}_{q};z] = \sum_{n=0}^{\infty} \frac{(a_{1})_{n}...(a_{p})_{n}}{(b_{1})_{n}...(b_{q})_{n}} \frac{z^{n}}{n!}$$
 (A.1)

where  $(a)_n = \Gamma(a+n)/\Gamma(a)$  stands for the Pochhammer symbol. We suppose that none of the denominator parameters is a negative integer or zero. This series converges for all values of z, real or complex, when  $p \leq q$ , and for |z| < 1 when p = q + 1. In the latter case, it also converges (absolutely) on the circle |z| = 1 if Re  $(\sum_{i=1}^q b_i - \sum_{i=1}^p a_i) > 0$ . If p > q + 1, the series never converges, except either when z = 0 or when the series terminates, that is when one at least of the a parameters is zero or a negative integer.

Hypergeometric series admit in general an integral representation of the Euler's type [10] [11] [12], which permits the corresponding analytic continuation in the complex z-plane beyond the unit disc.

For the ordinary hypergeometric series, we have:

$$_{2}F_{1}[a,b;c;z] = \frac{\Gamma(c)}{\Gamma(b)\Gamma(c-b)} \int_{0}^{1} du \ u^{b-1} (1-u)^{c-b-1} (1-zu)^{-a}$$
 (A.2)

with Re(c) > Re(b) > 0. In order to get a single-valued analytic function in the whole complex z-plane we will follow the customary convention of assuming a cut along the real axis from 1 to  $\infty$ .

For the generalized hypergeometric function of one variable the integral representation of the Euler's type is:

$${}_{p}F_{q}[a_{1},\{a\}_{p-1};b_{1},\{b\}_{q-1};z] = \frac{\Gamma(b_{1})}{\Gamma(a_{1})\Gamma(b_{1}-a_{1})} \int_{0}^{1} du \, u^{a_{1}-1} (1-u)^{b_{1}-a_{1}-1} \, {}_{p-1}F_{q-1}[\{a\}_{p-1};\{b\}_{q-1};zu]$$
(A.3)

under the constraints:  $p \le q+1$ ,  $Re(b_1) > Re(a_1) > 0$  and none of  $b_i$ , i = 1...q, is zero or a negative integer, giving the analytic continuation in the whole complex z-plane, cut along the positive axis from 1 to  $\infty$  again.

A well-known transformation between Gauss' hypergeometric functions of one variable, needed in the main text is: [12]

$${}_{2}F_{1}[a,b;c;z] = (1-z)^{c-b-a} {}_{2}F_{1}[c-a,c-b;c;z]$$

$$= (1-z)^{-a} {}_{2}F_{1}[a,c-b;c;z/(z-1)] = (1-z)^{-b} {}_{2}F_{1}[c-a,b;c;z/(z-1)]$$
(A.4)

An interesting relation between an ordinary Gauss' function and an elementary function not usually shown in specialized tables is:

$$z_{2}F_{1}[1,2;3;z] = -2\left[1 + \frac{1}{z}\ln(1-z)\right]$$
 (A.5)

which can be proved by expanding both sides as power series.

### The dilogarithm and its relation to the $_3F_2$ function

The dilogarithmic function is defined as: [13] [6]

$$Li_2(z) = -\int_0^1 du \, \frac{\ln(1-zu)}{u}$$
 (A.6)

for values of z real or complex. If |z| < 1, the dilogarithm may be expanded as the power series:

$$Li_2(z) = \sum_{n=0}^{\infty} \frac{z^{n+1}}{(n+1)^2}$$
 (A.7)

corresponding to the principal value. We can also write:

$$Li_2(z) = z \int_0^1 du \ _2F_1[1, 1; 2; zu] = z \ _3F_2[1, 1, 1; 2, 2; z]$$
 (A.8)

The derivative of the dilogarithm is

$$\frac{d}{dz}Li_2(z) = -\frac{\ln(1-z)}{z} \tag{A.9}$$

## $\mathbf{B}$

Generalized Gauss' Functions of two variables: Appell's functions

In this paper we are involved in particular with the  $F_3$  Appell's function [5] [9] [10] [11], so we write its series expansion:

$$F_3[a, a', b, b'; c; x, y] = \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(a)_m (a')_n (b)_m (b')_n}{(c)_{m+n}} \frac{x^m}{m!} \frac{y^n}{n!}$$
(B.1)

which exists for all real or complex values of a, a', b, b', and c except c a negative integer. With regard to its convergence, the  $F_3$  series is absolutely convergent when both |x| < 1 and |y| < 1. Then there is no problem with internal rearrangements of the series.

The  $F_3$  function can be rewritten in terms of ordinary Gauss' functions:

$$F_3[a, a', b, b'; c; x, y] = \sum_{m=0}^{\infty} \frac{(a)_m(b)_m}{(c)_m} \frac{x^m}{m!} {}_{2}F_1[a', b'; c + m; y]$$
 (B.2)

where we have made use of the relation:  $(c)_{m+n} = (c)_m (c+m)_n$ .

Moreover, the  $F_3$  function admits the following integral representation: [5] [10]

$$F_3[a, a', b, b'; c; x, y] = \frac{\Gamma(c)}{\Gamma(b)\Gamma(b')\Gamma(c - b - b')} \times$$
(B.3)

$$\int \int du \ dv \ u^{b-1}v^{b'-1}(1-u-v)^{c-b-b'-1}(1-xu)^{-a}(1-yv)^{-a'}$$

where the integral is taken over the triangular region  $0 \le u$ ,  $0 \le v$ ,  $u + v \le 1$ , under the conditions: Re(b) > 0, Re(b') > 0, Re(c - b - b') > 0. This expression furnishes a single-valued analytic function in the domain defined by the Cartesian product of the complex planes of x and y with the restrictions  $|arg(1-x)| < \pi$ ,  $|arg(1-y)| < \pi$ . Hence, the order of integration may be reversed according to Fubini's theorem.

Some properties and relations between Appell's functions needed in this paper are given below: [5] [10]

$$F_1[a;b,b';c;x,x] = {}_2F_1[a,b+b';c;x] = {}_2F_1[b+b',a;c;x]$$
 (B.4)

$$F_3[a, c-a, b, b'; c; x, y/(y-1)] = F_3[b', b, c-a, a; c; y/(y-1), x] = (1-y)^{b'} F_1[a; b, b'; c; x, y]$$
(B.5)

$$F_3[a, c-a, b, c-b; c; x, y] = F_3[c-a, a, c-b, b; c; y, x] = (1-y)^{a+b-c} {}_2F_1[a, b; c; x+y-xy]$$
(B.6)

# The polylogarithm and its relation to the generalized Campé de Fériet function ${\cal F}_B^{(2)}$

The polylogarithm  $Li_q(z)$  is defined as a series as

$$Li_q(z) = \sum_{n=0}^{\infty} \frac{z^{n+1}}{(n+1)^q} \quad (q > 1)$$
 (B.7)

which can be expressed according to

$$Li_{q}(z) = z \int_{0}^{1} \ du \ _{q}F_{q-1}[\{1\}_{q};\{2\}_{q-1};zu] \ = \ z \ _{q+1}F_{q}[\{1\}_{q+1};\{2\}_{q};z]$$

A generalized Campé de Fériet function <sup>4</sup>, of particular interest for us, is defined as

$$F_B^{(2)}[\{b\}_r, \{b'\}_s; \{d\}_t; c; x, y] = \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(b_1)_m ... (b_r)_m (b'_1)_n ... (b'_s)_n}{(d_1)_m ... (d_t)_m (c)_{m+n}} \frac{x^m}{m!} \frac{y^n}{n!}$$
(B.8)

Thus the following equality is satisfied

$$x F_B^{(2)}[\{1\}_q, \{1\}_2; \{2\}_{q-2}; 3; x, y = 1] = 2 Li_q(x)$$
 (B.9)

which is the generalization of Eq. (5).

<sup>&</sup>lt;sup>4</sup>Campé de Fériet functions are special cases of generalized Lauricella functions of two variables [2]

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